

"Dark matter searches" with a focus on new techniques (Mono-X) WIN2013: Natal, Brazil

James D Pearce, On behalf of the ATLAS and CMS collaborations

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OUTLINE

- 1. Dark Matter Background
 - Evidence for Dark Matter
 - ► The "WIMP Miracle"
- 2. Effective Field Theories
- 3. Detection Methods
- 4. Mono-X Analyses
 - Monojet (ATLAS)
 - Monophoton (CMS)
 - ► Mono-W/Z (ATLAS + CMS)
 - ► Mono-b
- 5. Summary
- 6. Auxiliary Material

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EVIDENCE FOR DARK MATTER (I) Galactic Rotation Curves Strong



- Galactic rotation curves show stars orbit at the same speeds
- This implies mass density of galaxies is uniform.

Strong Gravitational Lensing



- Image of Abell 1689 cluster as observed by the hubble telescope
- The mass of galaxies is not enough to account for the strong gravitational lensing.

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EVIDENCE FOR DARK MATTER (II) Weak Gravitational Lensing Cosmic Microwave Background



- Two galaxy clusters colliding.
- The pink shows the x-ray emissions.
- Blue shows unseen mass as measured with weak gravitational lensing techniques.



- Anisotropies in the CMB are due to acoustic oscillations in the early universe.
- Angular scales of the oscillations reveal the different effects of baryonic matter and DM.

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RELIC ABUNDANCE AND THE "WIMP MIRACLE"



- 1. DM and SM particles are in thermal (chemical) equilibrium.
- 2. Universe expands and cools; DM production drops exponentially ($\sim e^{-m_{\chi}/T}$).
- 3. Energy drops below DM production threshold; DM abundance remains constant ("Freeze out").

We are left with a relic abundance of DM:

$$\Omega_\chi \propto rac{1}{\langle \sigma v \rangle} \sim rac{m_\chi^2}{g_{\chi^4}}$$

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WHAT WE KNOW ABOUT DARK MATTER

- 1. It's neutral under electric charge, since it does not produce photons,
- 2. It's stable, or at least has a lifetime on cosmological scales,
- 3. It's non-baryonic, to preserve the success of Λ CDM,
- 4. It has a relic abundance consistent with weak scale mass and interactions.

These seem to all point us to some sort of weakly interacting massive particle (WIMP). We can use an EFT to model what we know about DM, without resorting to any one specific UV complete theory (eg. SUSY, LED, etc.)

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FROM UV COMPLETE TO EFT (I)



By Taylor expanding the SM-DM propagator around the momentum transfer and only keeping the leading order we get an effective coupling constant:

$$rac{1}{Q_{tr}^2-M^2}=-rac{1}{M^2}\left(1+rac{Q_{tr}^2}{M^2}+\mathcal{O}\left(rac{Q_{tr}^4}{M^4}
ight)
ight)pprox -rac{1}{M^2}$$

This approximation is only valid if $Q_{tr}^2 \ll M^2$ otherwise all other terms in the expansion (UV complete theory) must be considered.

FROM UV COMPLETE TO EFT (II)

Once the mediator has been "integrated out" we no longer talk about the parameter M, instead we replace it with M_* , which parameterizes the energy scale of the EFT. M_* is the most important parameter of the theory, it's related to the mediator mass and couplings, and tells us where the EFT approach breaks down:

- $M_* = M/\sqrt{g_{\chi}g_q}$, where g_{χ} and g_q are the couplings of the mediator to the DM and quark fields.
- ► 4-momentum conservation requires $m_{\chi} < M/2$
- Perturbation theory requires $g_{\chi}g_q < (4\pi)^2$
- Therefore our EFT is valid for $m_{\chi} < 2\pi M_*$

This EFT language allows us to relate different experimental signatures in a model-independent way. As we'll see the relic abundance, direct detection signal, and collider predictions depend only on M_* .

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DETECTION

EFT'S

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BACKGROUND

INTRO

om UV	COMPLETE TO EFT (III)	alv.10

Monojet

MONOPHOTON

$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	Name	Operator	Coefficient	The theory is then
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	D1	$\bar{\chi}\chi\bar{q}q$	m_q/M_*^3	characterized by an effective
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	D2	$\bar{\chi}\gamma^5\chi\bar{q}q$	im_q/M_*^3	Lagrangian \mathcal{L}_{eff} :
$ \begin{array}{c ccccc} D4 & \bar{\chi}\gamma^5\chi\bar{q}\gamma^5q & m_q/M_*^3 \\ D5 & \bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q & 1/M_*^2 \\ D6 & \bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q & 1/M_*^2 \\ D7 & \bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma^5q & 1/M_*^2 \\ D8 & \bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q & 1/M_*^2 \\ D9 & \bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q & 1/M_*^2 \\ D10 & \bar{\chi}\sigma_{\mu\nu}\gamma^5\chi\bar{q}\sigma_{\alpha\beta}q & i/M_*^2 \\ D11 & \bar{\chi}\chi G_{\mu\nu}G^{\mu\nu} & \alpha_s/4M_*^3 \\ D13 & \bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu} & i\alpha_s/4M_*^3 \\ D14 & \bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu} & \alpha_s/4M_*^3 \\ \end{array} \right) \qquad $	D3	$ar{\chi}\chiar{q}\gamma^5 q$	im_q/M_*^3	$\mathcal{L}_{eff} = \sum c_i O_i$
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	D4	$ar{\chi}\gamma^5\chiar{q}\gamma^5q$	m_q/M_*^3	Where $c = \frac{1}{1}$ and O is an
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	D5	$\bar{\chi}\gamma^{\mu}\chi\bar{q}\gamma_{\mu}q$	$1/M_{*}^{2}$	$M_*^{d-4} \text{ and } O_1 \text{ is an } M_*^{d-4}$
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	D6	$\bar{\chi}\gamma^{\mu}\gamma^{5}\chi\bar{q}\gamma_{\mu}q$	$1/M_{*}^{2}$	effective operator which is
$ \begin{array}{c cccc} D8 & \bar{\chi}\gamma^{\mu}\gamma^{5}\chi\bar{q}\gamma_{\mu}\gamma^{5}q & 1/M_{*}^{2} \\ D9 & \bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q & 1/M_{*}^{2} \\ D10 & \bar{\chi}\sigma_{\mu\nu}\gamma^{5}\chi\bar{q}\sigma_{\alpha\beta}q & i/M_{*}^{2} \\ D11 & \bar{\chi}\chi G_{\mu\nu}G^{\mu\nu} & \alpha_{s}/4M_{*}^{3} \\ D12 & \bar{\chi}\gamma^{5}\chi G_{\mu\nu}G^{\mu\nu} & i\alpha_{s}/4M_{*}^{3} \\ D13 & \bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu} & i\alpha_{s}/4M_{*}^{3} \\ D14 & \bar{\chi}\gamma^{5}\chi G_{\mu\nu}\tilde{G}^{\mu\nu} & \alpha_{s}/4M_{*}^{3} \end{array} \right) \\ \begin{array}{c} \text{combination of the SM and} \\ DM (\chi) \text{ fields.} \\ P \text{lace limits on a} \\ \text{representative set: D1 (scalar),} \\ D5 (vector), D8 (axial-vector), \\ D9 (tensor) \text{ and D11 (couples)} \\ \text{to gluons)} \end{array} $	D7	$\bar{\chi}\gamma^{\mu}\chi\bar{q}\gamma_{\mu}\gamma^{5}q$	$1/M_{*}^{2}$	some Lorentz invariant
$ \begin{array}{c ccccc} D9 & \bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q & 1/M_*^2 \\ D10 & \bar{\chi}\sigma_{\mu\nu}\gamma^5\chi\bar{q}\sigma_{\alpha\beta}q & i/M_*^2 \\ D11 & \bar{\chi}\chi G_{\mu\nu}G^{\mu\nu} & \alpha_s/4M_*^3 \\ D12 & \bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu} & i\alpha_s/4M_*^3 \\ D13 & \bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu} & i\alpha_s/4M_*^3 \\ D14 & \bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu} & \alpha_s/4M_*^3 \\ \end{array} \right) \qquad $	D8	$\bar{\chi}\gamma^{\mu}\gamma^{5}\chi\bar{q}\gamma_{\mu}\gamma^{5}q$	$1/M_{*}^{2}$	combination of the SM and $DM(x)$ (11)
$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	D9	$\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q$	$1/M_{*}^{2}$	DM (χ) fields.
$ \begin{array}{c cccc} D11 & \bar{\chi}\chi G_{\mu\nu}G^{\mu\nu} & \alpha_s/4M_*^3 & \text{representative set: D1 (scalar),} \\ D12 & \bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu} & i\alpha_s/4M_*^3 & D5 (vector), D8 (axial-vector), \\ D13 & \bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu} & i\alpha_s/4M_*^3 & D9 (tensor) \text{ and D11 (couples to gluons)} \\ D14 & \bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu} & \alpha_s/4M_*^3 & to gluons) \end{array} $	D10	$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi\bar{q}\sigma_{\alpha\beta}q$	i/M_*^2	 Place limits on a
D12 $\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$ $i\alpha_s/4M_*^3$ D5 (vector), D8 (axial-vector), D13 $\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$ $i\alpha_s/4M_*^3$ D9 (tensor) and D11 (couples D14 $\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$ $\alpha_s/4M_*^3$ to gluons)	D11	$\bar{\chi}\chi G_{\mu u}G^{\mu u}$	$lpha_s/4M_*^3$	representative set: D1 (scalar),
D13 $\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$ $i\alpha_s/4M_*^3$ D9 (tensor) and D11 (couples D14 $\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$ $\alpha_s/4M_*^3$ to gluons)	D12	$\bar{\chi}\gamma^5\chi G_{\mu u}G^{\mu u}$	$i lpha_s / 4 M_*^3$	D5 (vector), D8 (axial-vector),
D14 $\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$ $\alpha_s/4M_*^3$ to gluons)	D13	$\bar{\chi}\chi G_{\mu u}\tilde{G}^{\mu u}$	$ilpha_s/4M_*^3$	D9 (tensor) and D11 (couples
	D14	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$lpha_s/4M_*^3$	to gluons)

SUMMARY 0

Mono-b

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Mono-W/Z



Direct detection

DETECTION METHODS



Experiments:

- ► ATLAS
- ► CMS
- ► D0
- ► CDF

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Experiments:

- ► XENON100
- ► CDMS
- ► SIMPLE
- ▶ CoGent
- ► IceCube
- ► Picasso
- ► COUPP

Indirect detection

Experiments:

- ► Fermi-LAT
- ► PAMELA
- ► AMS-02
- ► WMAP
- Planck

...and many more

MONOJET ANALYSIS (ATLAS/CMS)

MONOIET



EW background estimate (ATLAS):

MONO-W/Z

Mono-h

SUMMARY

•
$$N_{SR}^{est} = (N_{CR}^{Data} - N_{CR}^{bkg}) \times (1 - F_{EW}) \times TF$$

• where $1 - F_{EW} = \frac{N_{CR}^{MC}}{\sum_{CR}^{AII EW} N_{CR}^{MC}}$

• and
$$TF = \frac{N_{SR}^{MC}}{N_{CR}^{MC}}$$

Main Backgrounds:

INTRO

BACKGROUND

- ► $Z(\rightarrow \nu \nu)$ + jet(s) (50-70%)
- $W(\rightarrow l^{\text{inv}}\nu) + \text{jet(s)} (46-29\%)$
- ► $Z(\rightarrow l^{inv}l^{inv}) + jet(s)$ (4-0%)

CMS-PAS-EXO-12-048

MC EW background estimate (CMS):

- Z+ jets, W+ jets, $t\bar{t}$ and single top:
- MadGraph → Phythia6: Z2Star tune with CTEQ6L1 pdf

ATLAS-CONF-2012-147

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MONOJET ANALYSIS (CMS)



Selection

- Trigger: $E_T^{miss} > 80 \text{ GeV}$
- ► Lead jet: *p_T* > 110 GeV, |η| < 2.4</p>
- lepton veto: e, μ , τ
- ► $\Delta \phi(\text{jet}_1, \text{jet}_2) < 2.5$
- ► jet veto: $N_{\text{jet}} \leq 2$
- Scan in E_T^{miss}

Blue dashed line indicates hypothetical DM signal



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ENERGY SCALE LIMITS



- Green line indicates the M_{*} values at which WIMPs of a given mass would result in the required relic abundance.
- M* limits above the thermal relic line means exclusion or negative interference or additional annihilation (e.g. to leptons)
- ► The light-grey region indicates where the EFT breaks down.

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WIMP-NUCLEON SCATTERING LIMITS



- Cross sections above observed are excluded.
- Assumption is that DM interacts with SM particles solely by a given operator: SI = D5, SD = D8
- Yellow contours show candidate events from CDMS: arXiv:1304.4279

WIMP ANNIHILATION LIMITS

MONOIET

INTRO

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 Comparison with FERMI-LAT is possible through our EFT The results can also be interpreted in terms of limits on WIMPs annihilating to light quarks

Mono-h

SUMMARY

MONO-W/Z

- All limits shown here assume 100% branching fractions of WIMPs annihilating to quarks
- Below 10 GeV for D5 and 70 GeV for D8 the ATLAS limits are below the values needed for WIMPs to make up the DM relic abundance

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MONOPHOTON ANALYSIS (CMS)



Main Backgrounds:

- $Z\gamma \rightarrow \nu\nu + \gamma$ (60%)
- $W\gamma \rightarrow l^{\text{inv}}\nu + \gamma$, di- γ and γ + jet (5%)
- ► Fake photons (20%) CMS-EXO-11-096





- All backgrounds estimated with MC
- ► Z/γ (NLO), di- γ and γ + jet ← Pythia6 with CTEQ6L1
- $W\gamma(NLO) \leftarrow MadGraph5$



CMS Experiment at LHC, CERN Data recorded: Sun Apr 24 22:57:52 2011 CD RuntEvent: 165374 / 314736281 Lumi section: 604

CERN-PH-EP-2012-209

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MONOPHOTON ANALYSIS

Source	Estimate
Jet Mimics Photon	11.2 ± 2.8
Beam Halo	11.1 ± 5.6
Electron Mimics Photon	3.5 ± 1.5
$W\gamma$	3.0 ± 1.0
γ +jet	0.5 ± 0.2
$\gamma\gamma$	0.6 ± 0.3
$Z(\nu \bar{\nu})\gamma$	45.3 ± 6.9
Total Background	75.1 ± 9.5
Total Observed Candidates	73



Selection:

- ► Trigger: Single photon
- ► Photon: *p_T* > 145 GeV, |η| < 1.44 (barrel region)</p>
- Energy ratio: HCAL/ECAL < 0.05 within $\Delta R < 0.15$
- Lepton veto and hadronic activity veto
- Isolated photons
- ► jet veto: $N_{\text{jet}} \leq 2$
- SR: $E_T^{miss} > 130 \text{ GeV}, |\eta| < 4.5$



WIMP-NUCLEON SCATTERING LIMITS



- Cross sections above observed are excluded.
- Spin-dependent/independent limits correspond to D8 and D5 operators.
- ► Not sensitive to D11 (gluon) operator.

Mono-W/Z Analyses (ATLAS/CMS)



INTRO

BACKGROUND

Two possible diagrams lead to interference:

- 1. $\xi = -1$: signal enhanced
 - leads to stronger signal than monojet!
- 2. $\xi = +1$: signal suppressed

Two different strategies:

 ATLAS: look for a single fat-jet with internal structure (mono-W/Z)

MONO-W/Z

 Same backgrounds as monojet analysis

Mono-h

SUMMARY

- Similar data-driven background estimate
- 2. CMS: look for single lepton (mono-W)
 - $W \rightarrow l\nu, t\bar{t}$, single top, Drell-Yan, diboson
 - Background estimated from MC

ATLAS-CONF-2013-073

CMS-EXO-12-060

MONO-W/Z A.K.A. MONO-FATJET (ATLAS)

Process	$E_{\rm T}^{\rm miss} > 350 {\rm GeV}$	$E_{\rm T}^{\rm miss} > 500 {\rm GeV}$
$Z \rightarrow \nu \bar{\nu}$	400+39	54^{+8}_{-10}
$W \to \ell^\pm \nu, Z \to \ell^\pm \ell^\mp$	210+20	22+4
WW, WZ, ZZ	57 ⁺¹¹	$9.1^{+1.3}_{-1.1}$
tī, single t	39 ⁺¹⁰	$3.7^{+1.7}_{-1.3}$
Total	710+48 -38	89 ⁺⁹ -12
Data	705	89



Selection:



SUMMARY

• Trigger: $E_T^{miss} > 80 \text{ GeV}$

MONO-W/Z

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MONO-h

- ► Fat jet: Cambridge-Aachen algorithm, R = 1.2, first two sub-jets balanced ($\sqrt{y} > 0.4$), $p_T > 250 \text{ GeV}, |\eta| < 1.2,$ $m_{jet} = 50 - 120 \, \text{GeV}$
- Veto leptons (e, μ) and photons
- $t\bar{t}$ supression: veto if >= 2 AntiKt4 jets with $\Delta R(\text{jet}_{Fat}, \text{jet}_{AntiKt4}) > 0.9 \text{ OR}$ $\Delta \phi(E_T^{miss}, \text{jet}) < 0.4$ for any AntiKt4 jets.
- SR: $E_T^{miss} > \{350, 500\}$ GeV

INTRO

BACKGROUND

MONO-W A.K.A. MONO-LEPTON (CMS)

MONOIET



INTRO

BACKGROUND

Selection:

 Trigger: single muon (*p_T* > 40 GeV) or electron (*p_T* > 80 GeV) triggers

MONO-h

MONO-W/Z

- Muons: $p_T > 45$ (offline) GeV, $|\eta| < 2.1$, isolated
- Electrons: $p_T > 100$ (offline) GeV, $|\eta| < 1.442$ or $1.566 < |\eta| < 2.5$, isolated
- ► Back-to-back kinematics: $0.4 < p_T / E_T^{miss} < 1.5 \text{ AND}$ $\Delta \phi(E_T^{miss}, l) > 0.8\pi$
- SR: $m_T > \{1, 1.5, 2\}$ TeV



SUMMARY

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WIMP-NUCLEON SCATTERING LIMITS



- Cross sections above observed are excluded.
- Assuming constructive interference gives better limits then ATLAS monojet and monophoton combined!
- ► Not sensitive to D11 (gluon) operator.

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Mono-b



Motivation:

- ► D1 is proportional to the initial quark mass (D1 ~ $\frac{m_q}{M_*^2}$)
- By approximate QCD flavour symmetry the outgoing quark will be a b
- Analysis for 2012 data is currently underway (ATLAS)

arXiv:1307.7834

 Despite the kinematic and PDF suppression for producing third generation quarks improvement on limits is up to 3 orders of magnitude!



INTRO BACKGROUND EFT'S DETECTION MONOJET MONOPHOTON MONO-W/Z MONO-b SUMMARY 0000 000 0 0000 000 0000 0 •

SUMMARY

- WIMPs are well motivated by what we know about Dark Matter and the observed relic abundance.
- EFT's allow us to search for WIMPs in a model-independent way as well as compare results from different experiments and signatures.
- Mono-X searches at the LHC are competitive and complementary to direct and indirect detection experiments.

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Auxiliary slides

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MONOJET ANALYSIS(ATLAS)



 Orange dashed line indicates hypothetical DM signal (×5) Selection:

- Trigger: $E_T^{miss} > 80 \text{ GeV}$
- At least one primary vetex
- ► Lead jet: $p_T > 120$ GeV, $|\eta| < 2$
- lepton veto: e, μ
- Multijet suppression: $\Delta \phi(E_T^{miss}, \text{jet}_2) > 0.5$
- ► jet veto: $N_{\text{jet}} \leq 2$

► SR:

 $E_T^{miss} > \{120, 220, 350, 500\}$ GeV jet $p_T > \{120, 220, 350, 500\}$ GeV
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WIMP-NUCLEON SCATTERING LIMITS

arXiv:1210.4491



- Cross sections above observed are excluded.
- Assumption is that DM interacts with SM particles solely by a given operator

CMS-PAS-EXO-12-048 WIMP-NUCLEON SCATTERING LIMITS

INTRO

BACKGROUND



MONOPHOTON

MONO-W/Z

MONO-h

SUMMARY

- ► CMS (2012) limits for D5 (vector) operator
- Light mediator model is studied to see how limits change with mediator mass, WIMP mass and decay width.